

FAST AND ACCURATE PROPAGATION OF COHERENT LIGHT

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Abstract. We describe a fast algorithm to propagate, for any user-specified accuracy, a time-harmonic electromagnetic field between two parallel planes separated by a linear, isotropic, and homogeneous medium. The analytic formulation of this problem (circa 1897) requires the evaluation of the so-called Rayleigh-Sommerfeld integral. If the distance between the planes is small, this integral can be accurately evaluated in the Fourier domain; if the distance is very large, it can be accurately approximated by asymptotic methods. In the large intermediate region of practical interest

$$u(\mathbf{x}, z) = \frac{1}{z} \int_{\mathbb{R}^2} f(\mathbf{y}) \frac{e^{i2\pi R}}{R} d\mathbf{y}, \quad z > 0,$$

where $R = \sqrt{z^2 + |\mathbf{x} - \mathbf{y}|^2}$. Given the input $u(\mathbf{y}, 0)$ in the plane $z = 0$, we seek the output $u(\mathbf{x}, z)$ for $z > 0$ that satisfies the free-space wave equation. The exact solution is given by the Rayleigh-Sommerfeld diffraction integral. In this work, we propose a fast and accurate propagation method based on the Fourier transform.

numerical simulation for propagation of a plane wave through a lens. The numerical simulation is based on the Rayleigh-Sommerfeld formula. The numerical simulation is based on the Rayleigh-Sommerfeld formula.

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Preliminaries

The Rayleigh-Sommerfeld Formula. The behavior of a wave function $\psi(\mathbf{r}, t)$ in a homogeneous medium is governed by the wave equation

Given the boundary data $u(x, z=0)$ and $f(x)$ we rewrite (1) as

$$u(x, z) = \int_{\mathbb{R}^2} f(y) K_z(x-y) dy,$$

where the kernel $K_z(r)$ is given by

$$K_z(r) = \frac{e^{i2z\sqrt{1+(r/z)^2}}}{iz} - \frac{1}{r/z} - \frac{i}{z\sqrt{1+(r/z)^2}}, \quad r \geq 0.$$

Denote the Fourier transform of the boundary data as

$$f(p) = \int_{\mathbb{R}^2} f(x) e^{-i2x \cdot p} dx,$$

we write (1) in the Fourier domain as

$$u(x, z) = \int_{\mathbb{R}^2} f(p) K_z(p) e^{i2x \cdot p} dp,$$

where the Fourier transform of the kernel $K_z(r)$ is an entire function satisfying

$$K_z(p) = e^{i2z\sqrt{1-p^2}}, \quad |p| \leq 1/z.$$

Our goal is to validate a numerical scheme that the computational cost does not increase with the span z . It is clear that the spatial kernel $K_z(r)$ is a high oscillatory function of r when z is small and that the Fourier transform $K_z(p)$ is a high oscillatory function of p when z is large. For an physical interpretation, note that the span z in the input data $u(x, z=0)$ and $K_z(p)$ are both high oscillatory functions of p and r on $K_z(r)$ and $K_z(p)$ respectively. In § 2 we will show how to approximate the convolution error and then describe a numerical method to approximate the solution using Green's function to boundary data $u(x, z=0)$ and the analytical series solution propagation problem or alternatively at an arbitrary value of z . For small values of z it is well known that the problem is a boundary value problem in the complex plane. For large values of z the problem is a boundary value problem in the complex plane. § 2.1. A and B. Both on the support.

Remark 1. Given the non-regularity of the boundary data

$$\frac{u(x, z)}{z} = g(x), \quad z=0,$$

for a high order uniform approximation is

$$u(x, z) = \int_{\mathbb{R}^2} g(y) \frac{e^{i2R}}{R} dy, \quad R = z^2 + |x-y|^2, \quad z > 0.$$

The numerical solutions our approach is also applicable to validation.

Slepian Functions. A phs ra st ust vntua a n spa an at th sa t ar ss nt a ban t nth Four r o an An appr at at a s rpt on o su h is was nt at b p an an h s o labor ators n b ons rn a spa t nan ban t n ra op rator an us n ts n un t ons to nt a ass o un t ons that hav ontro on ntr at on n both th spa an th Four r o ans p an t a show that th s nt ra op rator o ut s w th r nt a op rator o ass a at a ph s s r b n th pro at sp h ro a w av un t ons both op rators shar th sa n un t ons

For our purposes we use n un t ons w th ontro on ntr at on n a squar e th spa t a o an an ban t to a s n th Four r o an h

The Unequally Spaced Fast Fourier Transform. In order to evaluate the forward transform

$$y_{mm'} = \sum_{m,m'=1}^M f_{mm'} e^{ix_{mm'}}$$

at output points $x_{nn'}$ where $n, n' = 1, \dots, N$ is usually done by using the FFT algorithm with computational complexity $\mathcal{O}(N^2)$.

in z is much larger than the spatial extent of $f(\mathbf{y})$ it is often to a further approximation $e^{i\frac{\pi}{z}\|\mathbf{y}\|^2}$ which when used in the Fraunhofer so that is a further approximation

$$\mathbf{u}(\mathbf{x}, z) = \frac{e^{i2\pi z} e^{i\frac{\pi}{z}\|\mathbf{x}\|^2}}{iz} \mathbf{f}\left(\frac{\mathbf{x}}{z}\right)$$

is the Fraunhofer approximation which is the output to the Fourier transform of the input spectrum of an antenna system. The ratio of the area of the aperture to the area of the observation plane is

$$\mathbf{u}(\mathbf{x}, z) = -\frac{ize^{i2\pi} \frac{z^2 + \|\mathbf{x}\|^2}{z^2}}{\mathbf{x}^2} \mathbf{f}\left(\frac{\mathbf{x}}{z^2 + \|\mathbf{x}\|^2}\right), \quad z \rightarrow \infty$$

which is a better approximation to the Fraunhofer approximation in the region of support.

A New Algorithm for Fast and Accurate Light Propagation

In this section we describe a fast algorithm to compute the output of a system z and an input spectrum $f(\mathbf{y})$ for $z > 0$. $\mathbf{u}(\mathbf{x}, z)$ is a square output when $\mathbf{W} = \left[-\frac{w}{2}, \frac{w}{2}\right]^2$ assuming that the boundary of f has a finite support \mathbf{f} is a function of \mathbf{y} with support $\mathbf{y} \in \mathbf{B}$. Hence \mathbf{f} is a function of \mathbf{y} with support $\mathbf{y} \in \mathbf{B}$ and \mathbf{c} is a function of \mathbf{y} in a square aperture $\mathbf{A} = \left[-\frac{a}{2}, \frac{a}{2}\right]^2$ so that a or n to w in to output

$$\mathbf{u}(\mathbf{x}, z) = \mathbf{f}(\mathbf{y}) \mathbf{K}_z(\mathbf{x} - \mathbf{K}(\mathbf{m}))$$

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phas that n. th sr a ura κ s s a) b th propa at on stan z sn th a ntu o th rn.) a s.) z^{-1} a on th opt a)ax s. Insp r b th Fr sn.)approx at on w r wrt th rn.)as

$$K_z(r) = \frac{e^{i2z} e^{i\frac{\pi}{z}r^2}}{iz} A_z(r),$$

wh r

$$A_z(r) = \frac{1}{\sqrt{r/z}} \frac{i}{z \sqrt{r/z}} e^{i2z \left(\sqrt{1+(r/z)^2} - 1 - \frac{1}{2}(r/z)^2 \right)}.$$

where \mathbf{c} is the band vector to the input unit on \mathbf{f} in the band vector \mathbf{c}' with respect to the input rays or a similar expression using the quadrature form or

Let $\mathbf{y}_{mm'} = \mathbf{y}_m \cdot \mathbf{y}_{m'}$, $\mathbf{A}_{m,m'} = \dots, \mathbf{M}$ be the $\mathbf{M} \times \mathbf{M}$ tensor product of quadrature nodes with the corresponding quadrature weights $w_{m,m'}$ set an $\mathbf{N} \times \mathbf{N}$ matrix of output locations $\mathbf{x}_{nn'} = \mathbf{x}_n \cdot \mathbf{x}_{n'}$, $\mathbf{W}_{n,n'} = \dots, \mathbf{N}$ then apply the quadrature form or to the input rays and obtain an approximation to the output field at the source locations as

$$\mathbf{u}_{nn'} = \frac{e^{i2z} L}{iz} \sum_{m,m'=1}^M w_{m,m'} \mathbf{T}_{nn'mm'}^{(\cdot)} \cdot \mathbf{y}_{mm'} e^{i2 \ell \mathbf{x}_{nn'} \cdot \mathbf{y}_{mm'}}.$$

In the $\mathbf{N} \times \mathbf{N} \times \mathbf{M} \times \mathbf{M}$ fourth order tensors $\mathbf{T}^{(\cdot)}$

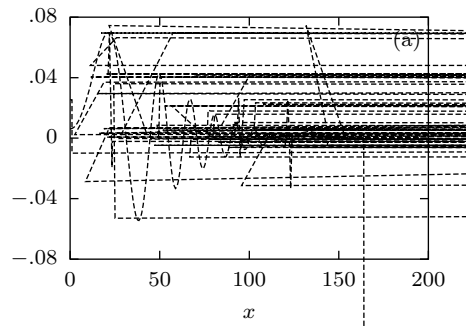
Lemma 4. Let $\mathbf{q}^{(l)}$, $\mathbf{U}_{nq}^{(l)}$, and $\mathbf{V}_{mq}^{(l)}$, where $l = 1, \dots, L$,

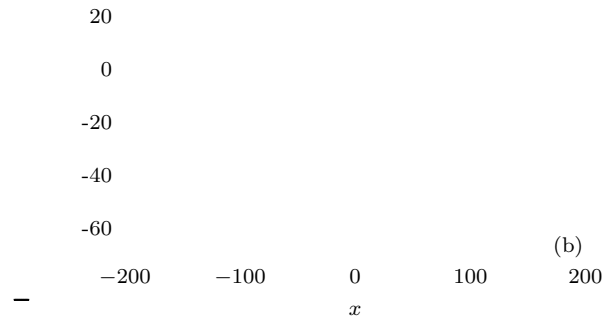
Theorem 5. The error of computing the field \mathbf{u} from $\mathbf{u}_{\text{nn}'}$ using $\mathbf{Q}_{\text{mm}'}$ is bounded by

$$|\mathbf{u}_{\text{nn}'}(z) - \mathbf{u}_{\text{nn}'}| \leq \frac{\mathbf{K} \mathbf{Q} \mathbf{R} \mathbf{f}_1}{z}.$$

The expression for $\mathbf{u}_{\text{nn}'}$ in (10) allows us to evaluate the error rapidly. First apply $\mathbf{Q}_{\text{mm}'}$ as a perturbation to the input signal $\mathbf{f}_{\text{mm}'}$.

to illustrate the relation between \mathbf{W}_{\max} and \mathbf{z}_{\min} or our theorem and \mathbf{W}'_{\max} and \mathbf{z}'_{\min} or the Fresnel approximation. It is obvious that $\mathbf{z}'_{\min} \approx \mathbf{z}_{\min} - 3 \lambda \mathbf{a}$ in the limit of large propagation distance $\mathbf{z} \gg \lambda$.





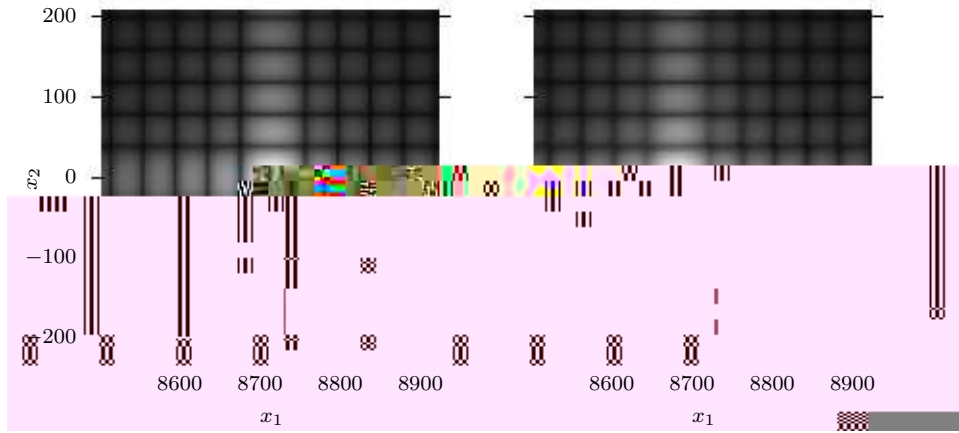


Figure 5.5. Comparison of the antenna of the Fourier transform point of the optical axis. The top part shows the intensity distribution of the antenna, and the bottom part shows the intensity distribution of the antenna. The x1 axis ranges from 8600 to 8900, and the x2 axis ranges from -200 to 200.

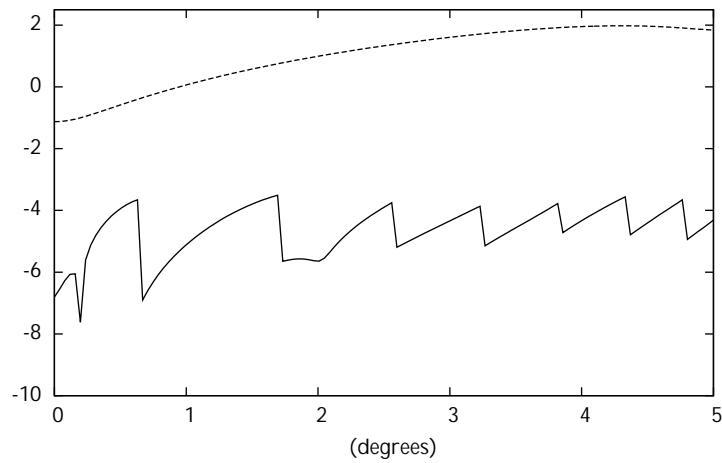


Figure 5.6. A graph showing the relationship between an angle in degrees (x-axis, 0 to 5) and a value (y-axis, -10 to 2). The solid line represents the antenna of the Fourier transform point, and the dashed line represents the antenna of the Fourier transform point.

in Fig. 1, to be treated in the present section, and as shown in our example, shows that the Fresnel approximation is not applicable to the spot in a transition to its position. The number of spots between the bottom and top of the bottom part of Fig. 1

Representative Examples of Computational Cost. In our previous work, the cost of our algorithm is given by the number of FFTs required. Let $\mathbf{R} = \mathbf{R}^{(1)} \dots \mathbf{R}^{(L)}$ where L is the number of steps to approximate the input. As it turns out, \mathbf{R} is a sparse matrix with \mathbf{z} which is the input in the approximation of the input. The number of operations is

[21] Lord Rayleigh,