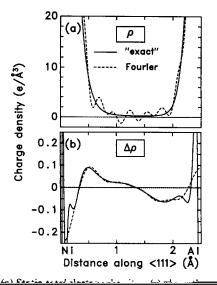
with the results of the self-consistent *ab initio* calculations based on the local density formalism [12, 13]. We focus on the following questions:

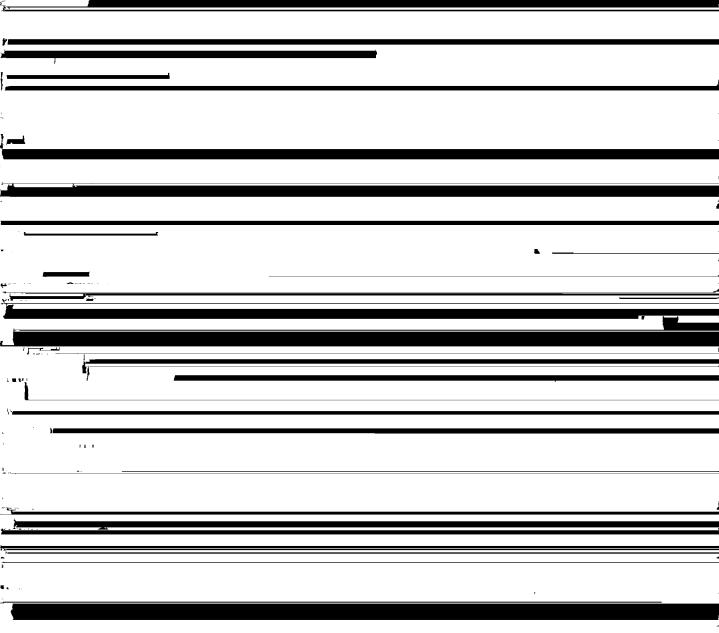
(i) How well can *ab initio* band theory describe the first few (low-angle) structure

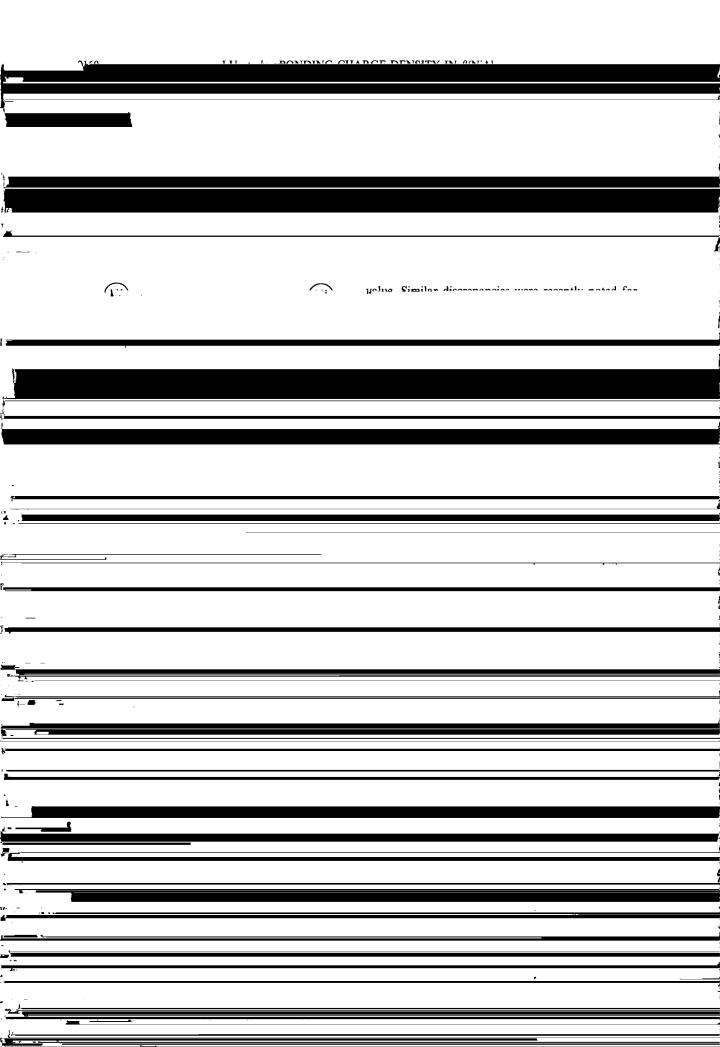
where the result naturally depends on the highest momentum  $(G_{\text{max}})$  included in this sum (as we will see below, current high-precision experiments are limited to rather small cut-off values  $G_{\text{max}}$ ). If the temperature factor can be deconvoluted from



representation the *total* density  $\rho$  [Fig. 1(a)] still exhibits significant oscillations in the bonding region. On the other hand, the density difference  $\Delta \rho_{\rm sup}({\bf r},G_{\rm big})$  (equation (12) and Fig. 1(b)] obtained from a Fourier series using  $G_{\rm big}=2\pi/a(6,3,1)$  closely mimics the directly calculated  $\Delta \rho_{\rm sup}({\bf r})$  in the bonding regions [of course,  $\Delta \rho_{\rm sup}({\bf r},G_{\rm big})$  still fails to reproduce the nodal structure near the core]. Note that the maximum magnitude of the static deformation density  $\Delta \rho_{\rm sup}({\bf r})$  outside the core is only  $\sim 0.1~{\rm e/\AA^3}$ , while the *total* density  $\rho({\bf r})$  has a magnitude of  $\sim 10~{\rm e/\AA^3}$  at this point. Clearly, the bonding charge is tiny.

Figure 2 shows as solid line the calculated static density deformation  $\Delta \rho_{\sup}(\mathbf{r})$  calculated without any Fourier truncation, comparing it to  $\Delta \rho_{\sup}(\mathbf{r}, G_{\text{big}})$  of equation (12), in which a large but finite cut-off  $G_{\text{big}} = 2\pi/a(6, 3, 1)$  (54 stars) was used. We see again that while the Fourier representation rounds off the





superposition atomic structure factors (using Hartree-Fock data [17] as an example) are E (100) = 12.20 and E (200) = 22.84 Hartes the

Here,  $B_{\rm Ni}$  and  $B_{\rm Al}$  are adjustable parameters while  $\rho_{\rm I}$  and  $\rho^{\rm MT}$  are fixed by theory (Table 1). This yields rather reasonable values of  $P_{\rm I}=0.56$  and  $P_{\rm I}=0.71$ 

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<u>.</u>	accurate measurements of crystalline structure factors to date. We also include in this table the "forbidden" (222) reflection measured by Alkire et al.	Static $\Delta \rho(\mathbf{r})$ Dynamic $\Delta F(\mathbf{r})$ $(100)+(110)+(111)+(200)$ $(100)+(110)+(111)+(200)$ $(100)+(110)+(111)+(200)$ $(100)+(110)+(111)+(200)$ $(100)+(110)+(111)+(200)$ $(100)+(110)+(111)+(200)$	
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- B. Dawson, Proc. R. Soc. A 298, 264 (1967); ibid 298, 379 (1967).
- 15. M. Deutsch, Phys. Lett. A 153, 368 (1991).
- J. E. Jaffe and A. Zunger, Phys. Rev. B 28, 5822 (1983).
- Z. W. Lu, S.-H. Wei, A. Zunger, S. Frota-Pessoa and L. G. Ferreira, *Phys. Rev. B* 44, 512 (1991).
- D. M. Ceperley and B. J. Alder, *Phys. Rev. Lett.* 45, 566 (1980).
- 27. J. P. Perdew and A. Zunger, *Phys. Rev. B* 23, 5048

(1968).

18. F. Herman and S. Skillman, Atomic Structure

- 28. S.-H. Wei and H. Krakauer, *Phys. Rev. Lett.* 55, 1200 (1985), and references cited therein.
- 20 D Village and I Calvert Daggeon's Handbook of

(1963). <u>p. 7. m. t</u> Crystallographic Data for Intermetallic Phases. Am.

3387 (1991).

- J. M. Zuo, J. C. Spence and M. O'Keeffe, *Phys. Rev. Lett.* 62, 2329 (1989).
- D. Hackenbracht and J. Kübler, J. Phys. F 10, 427 (1980).
   A. T. Hong and A. I. Freeman, Phys. Rev. B 43, 6446.
- P. Georgopoulos and J. B. Cohen, *Scripta metall.* 11, 147 (1977).
- 31. M. J. Cooper, Phil. Mag. 8, 811 (1963).
- 32. T. Hughes, E. P. Lautenschlager, J. B. Cohen and J. O. Brittain, J. appl. Phys. 42, 3705 (1971).
- 33 M Mosteller R M Nicklow D M Zehner S -C. Lui