



**PHYSICAL REVIEW E**

## II. MATERIALS, METHODS, AND TECHNIQUES

### A. Dispersion of gold nanoparticles in lyotropic LCs

GNRs of various sizes were examined. Comparatively large, polymer coated GNRs (Nanopartz Inc.) with a mean diameter of 40 nm and mean length of 73 nm are used as supplied. Smaller GNRs with a mean diameter of 20 nm and mean length of 50 nm were synthesized according to Ref. [29]. Gold nanocubes and gold “nanostars” were synthesized according to Refs. [30,31], respectively. Gold nanoparticles are functionalized by thiol-terminated methoxy-poly(ethylene glycol) (mPEG-SH) for colloidal stability. The composite of mPEG-GNRs and lyotropic LC is prepared based on a ternary lyotropic LC of sodium decyl sulfate-decanol-water (SDS–1-decanol–water) with a known phase diagram [32]. A calamitic nematic ( $N_C$ ) lyotropic LC was prepared using a composition of 37.5 wt% of SDS, 5.5 wt% of 1-decanol (both

FIG. 2. (Color online) (a) SEM image of gold nanocubes. (b) Optical microscopy image of a dispersion of gold nanocubes in 5CB. The inset shows a vial with the gold nanocubes-5CB dispersion. (c) Polarization-independent extinction spectrum of gold nanocubes in 5CB. (d)

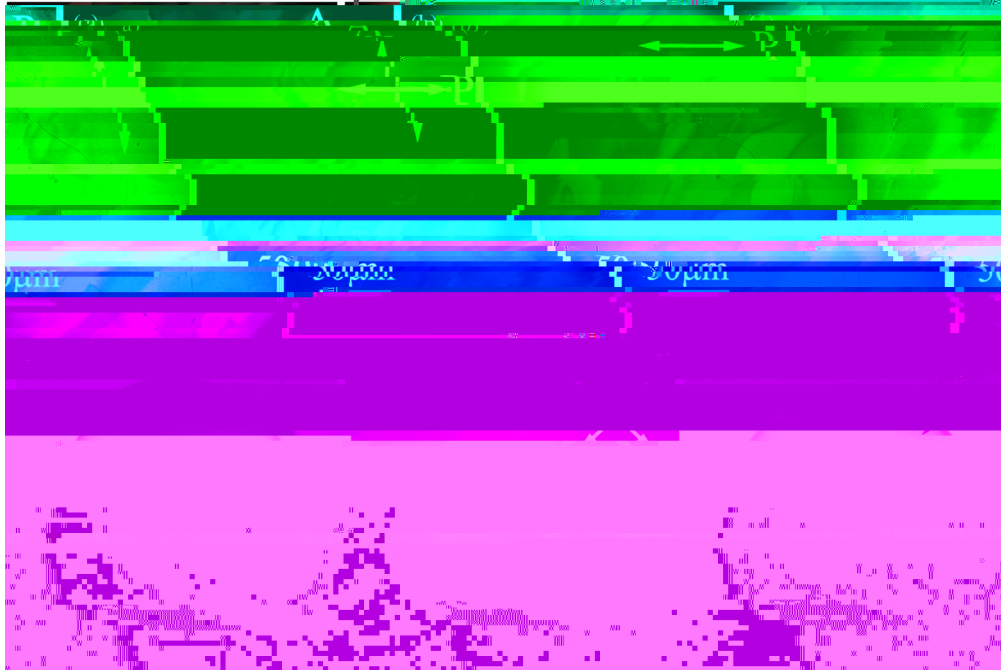


FIG 3. (a) POM image of 5CB liquid crystal phase. (b, d) Time evolution of 5CB liquid crystal phase in PDS cell. (c, e, f) Time evolution of 5CB liquid crystal phase in PDS cell.

Initial state  
 In the case of  
 big disturbance

$N(t) \sim N_0$ , as  
 indicated in  
 figure 4 (g)

5). Plot  
 $N$  and



FIG 4. (a) SEM image of the surface. (b, c) Time evolution of the system. (d, e) Time evolution of the system with a different parameter. (f, g) Time evolution of the system with a different parameter. (h) Plot of  $N$  and  $N_0$ .



FIG. 5. (a) A POM in a hybrid matrix 5CB

FIG. 6. (Color online) (a) A schematic of the setup used to measure the polarization-dependent transmitted spectra; the arrow indicates orientation of  $\mathbf{N}_0$  set by the shear force. (b) Experimental extinction spectra of GNRs at  $4.7 \times 10^{-8}$  M. (c) Simulated extinction coefficients for the same concentration of GNRs as (b) for two orthogonal polarizations of incident light and  $S_{\text{GNR}} = 1$ . (d) Computer-simulated spectral dependencies of extinction, absorption, and scattering coefficients for two orthogonal polarizations of incident light. (e) Calculated imaginary parts of refractive indices and its anisotropy. (f) Calculated spectral dependence of effective-medium refractive indices  $n$  and  $n$ .

at much lower concentrations or volume fractions of large GNRs. Third, the relative contribution of scattering to the total extinction ( $\frac{\text{abs}}{\text{ext}}$ ) increases with increasing the effective volume of GNRs. For example, this ratio of GNRs of 20 nm in diameter and 50-nm long is 12%, while it increases to 52% for GNRs of 40 nm in diameter and 73 nm in length. This property makes dispersions of large GNRs good candidates for applications based on light scattering.

The quality of orientational ordering of GNRs in LC can be characterized by the scalar order parameter defined as

the absorption coefficient  $\alpha_{\text{abs}}$  in Fig. 6(d):  $n_{\text{eff}}(\lambda) = n_{\text{offset}} = [1/(2\pi)] \text{P.V.} \int_{\lambda_1}^{\lambda_2} \alpha_{\text{abs}}(\lambda') / [1 - \tilde{S}(\lambda/\lambda')]^2 d\lambda'$ , where P.V. is the Cauchy principal value of the integral. The integration ranges from  $\lambda_1 = 450$  nm to  $\lambda_2 = 900$  nm. The used values of offset extraordinary and ordinary indices  $n_{\text{LC}}^{\text{offset}}$  are based on  $\bar{n}_{\text{LC}} = 1.39$  and the intrinsic optical anisotropy  $\Delta n_{\text{LC}} = 0.006 \pm 0.001$  of the  $\text{N}_C$ . The effective-medium optical anisotropy of the GNR- $\text{N}_C$  composite is much larger than the intrinsic birefringence of  $\text{N}_C$  and changes sign at around the longitudinal SPR peak wavelength  $\lambda_{\text{SPR}}$  [Fig. 6(f)].

#### IV. DISCUSSION

##### A. Strong anchoring and elastic alignment

In general, the alignment of rod-like or other nonspherical particles dispersed in the nematic LC could be caused by GNR-LC matrix interaction in “strong,” “weak,” or nite surface anchoring regimes. In the strong anchoring regime, the GNR-matrix interaction is expected to be mediated mostly by the minimization of elastic free energy due to director distortions induced by nanoparticles while having the director at LC-GNR surfaces follow the tangential boundary conditions [21]. In contrast, in the regime of weak surface anchoring, the director distortions in the LC bulk and energetic cost due to them can be neglected as the director meets the GNR-LC surface at different angles, so that the anisotropic nanoparticle orientation is determined by minimization of the

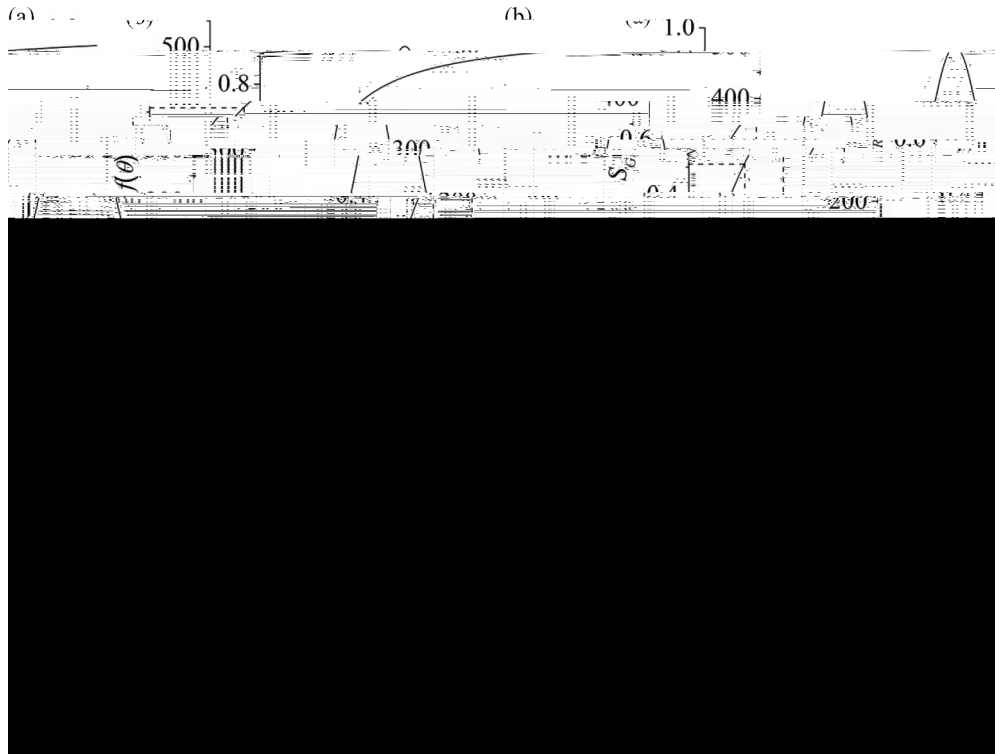


FIG. 7. (a) Plot of  $S_G$  vs  $r_f$  for  $S_G = 0.8$ . (b) Plot of  $S_G$  vs  $r_f$  for  $S_G = 1.0$ . (c) Plot of  $S_G$  vs  $r_f$  for  $S_G = 0.71$ . (d) Plot of  $S_G$  vs  $r_f$  for  $S_G = 0.5$ .

is a function of  $r_f$  and  $N$ . The function  $f_s(\theta)$  is defined as  $f_s(\theta) = \frac{1}{N} \sum_{i=1}^N \cos^2(\theta - \theta_i)$ , where  $\theta_i$  is the angle of the  $i$ -th particle. The function  $f_s(\theta)$  is a measure of the order of the system. The function  $f_s(\theta)$  is a function of  $\theta$  and  $N$ . The function  $f_s(\theta)$  is a function of  $\theta$  and  $N$ .

$$S_G = \int_0^{\pi} P_2(\theta) f_s(\theta) \sin \theta d\theta$$

$$= 3 \int_0^{\pi} \cos^2 \theta [2 - \cos^2(\theta - \theta_i)] \sin \theta d\theta$$



V. CONCLUSION

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