

Anisotropic electrostatic screening of charged colloids in nematic solvents

Jeffrey C. Everts^{1,2*}, Bohdan Senyuk^{3*}, Haridas Mundoor³, Miha Ravnik^{1,4†}, Ivan I. Smalyukh^{3,5,6†}

The physical behavior of anisotropic charged colloids is determined by their material dielectric anisotropy, affecting colloidal self-assembly, biological function, and even out-of-equilibrium behavior. However, little is known about anisotropic electrostatic screening, which underlies all electrostatic effective interactions in such soft or biological materials. In this work, we demonstrate anisotropic electrostatic screening for charged colloidal particles in a nematic electrolyte. We show that material anisotropy behaves markedly different from particle anisotropy. The electrostatic potential and pair interactions decay with an anisotropic Debye screening length, contrasting the constant screening length for isotropic electrolytes. Charged dumpling-shaped near-spherical colloidal particles in a nematic medium are used as an experimental model system to explore the effects of anisotropic screening, demonstrating competing anisotropic elastic and electrostatic effective pair interactions for colloidal surface charges tunable from neutral to high, yielding particle-separated metastable states. Generally, our work contributes to the understanding of electrostatic screening in nematic anisotropic media.

INTRODUCTION

Colloidal particles in nematic solvents exhibit a rich variety of self-assembly behaviors, ranging from isotropic to nematic, smectic, and cholesteric phases. The physical behavior of anisotropic charged colloids is determined by their material dielectric anisotropy, affecting colloidal self-assembly, biological function, and even out-of-equilibrium behavior. However, little is known about anisotropic electrostatic screening, which underlies all electrostatic effective interactions in such soft or biological materials. In this work, we demonstrate anisotropic electrostatic screening for charged colloidal particles in a nematic electrolyte. We show that material anisotropy behaves markedly different from particle anisotropy. The electrostatic potential and pair interactions decay with an anisotropic Debye screening length, contrasting the constant screening length for isotropic electrolytes. Charged dumpling-shaped near-spherical colloidal particles in a nematic medium are used as an experimental model system to explore the effects of anisotropic screening, demonstrating competing anisotropic elastic and electrostatic effective pair interactions for colloidal surface charges tunable from neutral to high, yielding particle-separated metastable states. Generally, our work contributes to the understanding of electrostatic screening in nematic anisotropic media.

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$$\phi(\mathbf{r}) = \mathcal{A}(\mathbf{r}, \lambda_D^I) \left(\frac{\lambda_D^I}{r} \right)^{\nu}, \quad (\text{Eq. 1})$$

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in the case of a function f of a variable x , we have

$$\frac{D(f)}{D(x)} = f'(x) + o(\epsilon) \quad (10)$$

where $\epsilon = \Delta x$ ($\Delta x > 0$), and $o(\epsilon)$ is a function of ϵ such that

... n_0 , ...
 ... (2) ...
 ...
 ... + ...
 ...
 ... (10) ...

... $\Phi(\dots)$, ...
 ... $\Phi(\dots)$, ...
 ... $\Phi(\dots)$

$$\Phi(\dots) = \Phi_E(\dots) + \Phi_{//}(\dots) + \Phi(\dots)$$

The first term in the above equation represents the energy of the system in the absence of the external magnetic field. The second term represents the energy of the system in the presence of the external magnetic field. The third term represents the energy of the system in the presence of the external magnetic field and the external electric field.

$$\Phi_E(\mathbf{r}, \mathbf{r}') = \alpha^2 \gamma \int \lambda_B(\mathbf{r}) \lambda_D(\mathbf{r}') \quad (1)$$

The above equation shows that the energy of the system is a function of the position of the particles. The energy of the system is a function of the position of the particles and the position of the external magnetic field. The energy of the system is a function of the position of the particles and the position of the external magnetic field and the external electric field.

... (51), ... (42, 52–54), ... (55), ... (56).

MATERIALS AND METHODS

Synthesis and characterization of charged colloidal dumpings

... (57), ... ()², ... ()², ... ()².

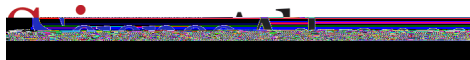
$\nabla^2 \phi(\mathbf{r}) = -\frac{1}{\epsilon_0} \rho(\mathbf{r})$, $\nabla \phi(\mathbf{r}) = -\mathbf{E}(\mathbf{r})$, $\nabla \times \mathbf{E}(\mathbf{r}) = 0$, $\nabla \cdot \mathbf{E}(\mathbf{r}) = \frac{\rho(\mathbf{r})}{\epsilon_0}$

$$\nabla^2 \phi(\mathbf{r}) = 0, \quad (\text{outside}) \quad (1)$$

\mathbf{n}

$$(\epsilon_1/\epsilon_2) \mathbf{n} \cdot \nabla \phi = \kappa^2 \phi$$

... (2) ...
... 2 ...
... (5)



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